Density functional studies of the electronic structure and adsorption at molybdenum oxide surfaces

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Abstract

The electronic structure and bonding at different oxygen sites of MoO 3 (010) and (100) surfaces is reviewed on the basis of ab initio density functional theory (DFT-LCGTO) cluster calculations. The clusters are chosen as finite sections of the ideal MoO 3 surface where cluster embedding is achieved by bond saturation with hydrogen terminator atoms yielding clusters up to Mo 7 O 30 H 18 . Resulting charge density distributions and binding properties are analyzed by populations, bond orders, and electrostatic potential maps. Interatomic binding at the surface is determined by both ionic and covalent contributions with a clear distinction between terminal oxygens and different bridging surface oxygens. Electronic differences between the MoO 3 (010) and (100) surfaces are found to be mainly due to the different atom arrangement while local atom charging and binding properties seem surface independent. The electronic surface parameters influence the behavior and reactions of adsorbed molecules as will be shown for H, OH, and C 3 H 5 adsorbates. © 1999 Elsevier Science B.V. All rights reserved.

Keywords: Molybdenum oxides; Binding; Adsorption; Reaction; Density functional theory; Cluster models

1. Introduction

Molybdenum oxides are of major scientific importance due to their complex chemical behavior. These oxides, prepared as pure substrates or as compounds including other transition metal oxides, can act as catalysts in many reactions of rather different type such as redox processes, acid–base processes (e.g., isomerization, polymerization), hydrogenation and dehydrogenation, selective oxidation, oxidative conversion [1]. In particular, molybdenum trioxide, MoO 3 , based materials become important as catalysts in the selective oxidation of hydrocarbons due to their structural and electronic surface properties. In experiments the oxidation is found to involve several steps yielding different products as a function of the local surface geometry and elemental composition of the molybdenum compound [1]. In a first step an organic molecule is activated by hydrogen abstraction near oxygen surface sites and in subsequent steps oxygens (possibly from different sites) are inserted. The details of these processes depend on the specific surface orientation. As an example we mention the allyl to acrolein conversion in a surface reaction

\[
\text{propene}(\text{CH}_3-\text{CH}-\text{CH}_2) \Rightarrow \text{allyl}(\text{CH}_2-\text{CH}-\text{CH}_2) \Rightarrow \text{acrolein}(\text{CH}_2=\text{CH}-\text{CHO})
\]

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PII: S0920-5861(98)00491-X
which involves simultaneous hydrogen abstraction and oxygen insertion and is found to occur on MoO$_3$(0 1 0) but not on (1 0 0) surfaces [1]. Experimental as well as theoretical details of these oxidation reactions are still under discussion [1–3] and a microscopic understanding of the underlying electronic structure and bonding properties at different surface oxygen sites is missing.

In the present theoretical study we examine the electronic structure and bonding near geometrically different surface oxygen sites for both (0 1 0) and (1 0 0) oriented ideal MoO$_3$ surfaces [4,5]. Since oxygen binding can be considered as a local phenomenon the surface cluster approach becomes meaningful. Thus, environments about the surface oxygens can be represented reasonably by bond saturated surface clusters of different size. Electronic wave functions and derived properties such as local charging, bond character, or electrostatic potentials of the surface clusters are determined by ab initio density functional theory (DFT) calculations. The results show that the interatomic binding at the two MoO$_3$ surfaces is determined by both ionic and covalent contributions with clear distinctions between terminal oxygens and the different bridging surface oxygens. Further, the difference in behavior between the (0 1 0) and (1 0 0) surface is found to be mainly due to the different atom arrangement while local atom charging and binding properties seem surface independent. The electronic structure at the MoO$_3$ surfaces influences possible reactions involving adsorbed molecules as will be discussed for hydrogen adsorption [6], OH desorption [6], and initial steps in the surface catalyzed allyl to acrolein oxidation [5,7].

2. Theoretical

Molybdenum trioxide, MoO$_3$, forms an ionic layer type orthorhombic crystal structure [8], where bilayers (described by (MoO$_3$)$_4$ unit cells) lie parallel to the (0 1 0) netplane which represents the easy cleavage plane of the crystal, see Fig. 1. While the internal interaction between atoms within the bilayers is dominated by rather strong ionic and covalent bonding the bilayers couple via weak (van der Waals) polarization forces. The ideal (0 1 0) surface of MoO$_3$ is characterized by a simple network of Mo and O ions where there are basically three types of structurally different surface oxygen centers, see Fig. 2(a). First, terminal (molybdenyl) oxygens, denoted O(a), are coordinated to one Mo center directly below and cover all Mo ions at the surface. Second, asymmetrically bridging oxygens, denoted O(b), are coordinated to one Mo center and couple weakly with another surface Mo. Third, symmetrically bridging oxygens, O(b'), are coordinated to two Mo centers of the surface and couple weakly with a Mo center of the underlying sublayer. The ideal (1 0 0) surface of MoO$_3$ is rather different in

![Fig. 1. Perspective view of the orthorhombic MoO$_3$ crystal lattice. The (0 1 0) and (1 0 0) netplanes as well as the Mo and O centers are indicated.](image-url)

![Fig. 2. Clusters used to model the MoO$_3$(0 1 0) surface: (a) defines the inequivalent oxygen centers O((a)–(d)) and lists interatomic distances (in Bohr); ((b)–(d)) show the clusters Mo$_7$O$_{30}$H$_{18}$, Mo$_3$O$_{16}$H$_{14}$, and Mo$_2$O$_{11}$H$_{10}$, where saturating hydrogens are sketched as small white spheres.](image-url)
its structure from the (0 1 0) surface, see Fig. 1. This surface is much less compact and there are many structurally different metal as well as oxygen sites, see Fig. 3(a). Molybdenyl oxygens, O(a), are similar to those of the (0 1 0) surface and coordinated to one Mo center. However, they do not cover all Mo ions leaving bare metal sites at the (1 0 0) surface. Symmetrically bridging oxygens, O(b′), correspond structurally to O(b′) atoms of the (0 1 0) surface. They are coordinated to two Mo centers of the surface and couple weakly with a third Mo surface center. Finally, oxygens O(d1) and O(d2) correspond structurally to terminal molybdenyl oxygens of the (0 1 0) surface. They lie close to the plane of the surface molybdenums and are coordinated to one Mo center.

The MoO₃ crystal structure determines the geometry of local model clusters for surface studies. The MoO₃(0 1 0) surface is represented in the present work by different clusters, shown in Fig. 2, where the largest cluster, Mo₇O₃₀H₁₈ (Fig. 2(b)), forms a symmetric surface section about a molybdenyl unit such that at least one member of each of the structurally different surface oxygens, O(a), O(b), and O(b′), occurs and experiences its full nearest neighbor environment within the cluster. In addition, peripheral oxygens are electronically saturated by hydrogen atoms. The MoO₃(1 0 0) surface is approximated by clusters shown in Fig. 3. Here the largest cluster, Mo₈O₂₄H₁₂ (Fig. 3(b)), yields also a characteristic surface unit containing each of the structurally different surface oxygens, O(a), O(b′), and O(d), with hydrogens serving as bond saturators at the cluster periphery. All surface clusters are chosen as finite sections of the ideal unreconstructed MoO₃(0 1 0) and (1 0 0) surfaces, respectively, which assumes that these surfaces do not exhibit major reconstruction as suggested from the experiment [1,9].

The electronic wave functions and derived properties of the clusters are determined by the ab initio density functional theory (DFT) method [10], using linear combinations of atomic orbitals (LCGTO) with flexible basis sets of contracted Gaussians. The basis sets are taken from atom optimizations and are all-electron type for all oxygen and hydrogen centers, while molybdenum centers are represented by valence basis sets with their [Ar]3d¹⁰ cores described by model core potentials. In the calculations the program package DeMon was applied using the local spin density approximation (LSDA) as well as generalized gradient schemes (GGA) for exchange and correlation. For further details see Ref. [4].

3. Results and discussion

3.1. MoO₃(0 1 0) and (1 0 0) substrate

The calculated values of atomic charges from Mulliken population analyses and of Mayer bond orders for the clusters modeling the MoO₃(0 1 0) and (1 0 0) surfaces are listed in Tables 1 and 2, respectively. As an important result from the calculations one finds only small variations between the different size clusters for a given surface orientation. Amongst the (010) surface clusters, the largest changes, 0.2 e, occur for the symmetrically bridging oxygens, O(b′), and for the nearby molybdenums in going from the two smaller to the largest cluster, Mo₇O₃₀H₁₈, cp. Table 1. This can be understood by the fact that in Mo₇O₃₀H₁₈ the O(b′) centers experience their natural Mo neighbors, see Fig. 2, whereas in the smaller clusters one of the two metal atoms is simulated by a hydrogen terminator.

Fig. 3. Clusters used to model the MoO₃(1 0 0) surface: (a) defines the inequivalent oxygen centers O((a)–(d)) about the two central molybdenums Mo(1) and Mo(2) and lists interatomic distances (in Bohr); (b)–(d) show the clusters Mo₆O₂₄H₁₂, Mo₅O₁₉H₈, and Mo₃O₁₂H₆.
The charges $q$ and bond orders $P$ involving molybdenum refer to the central Mo centers of the clusters.

The population analyses confirm the ionic nature of the MoO$_3$ compound as suggested by basic chemical intuition. However, the actual ionic charging is found to be smaller than expected which suggests sizable covalent contributions to the interatomic binding in the clusters. The Mo metal centers become positively charged and are described by the populations as Mo$^{x+}$. Allogether, the oxygen centers of the MoO$_3$(1 0 0) based clusters differ in their populations very little from those of the MoO$_3$(0 1 0) based clusters if structurally equivalent centers are compared. This suggests that charging of the atom centers is determined by the bulk structure rather than by the specific orientation of the ideal surface.

Results from bond order analyses give further information about interatomic binding at the MoO$_3$ surfaces. The Mo–O(a) bond order indices involving terminal oxygens yield values 1.93–2.00 which describes a molybdenyl double bond. The asymmetric bridging oxygens O(b) yield Mo–O bond orders that are quite different for the two Mo centers being bridged. For Mo closest to the oxygen, values of 1.5–1.6 show strong binding while for Mo further away bond orders of 0.2 suggest only small coupling. This confirms that the oxygens O(b) are coordinated mainly to one molybdenum center, and therefore, resemble the terminal molybdenyl oxygens O(a) which may be expected already from simple geometric considerations. Oxygenes O(b') bridging two Mo centers symmetrically result in Mo–O bond orders of 0.7–0.9 which reflects Mo–O single bonds with each of the neighboring molybdenums.

The different types of bonding for different surface oxygens become also obvious by an inspection of the charge rearrangement due to bond formation as evidenced in respective charge density difference, $\Delta \rho(r)$, maps which are shown elsewhere [4]. Further, electrostatic potentials $V(r)$ computed from the cluster charge distribution [4] can give information about local charging and binding at the surface. Results of the electrostatic potential $V(z)$ above different oxygen sites on the MoO$_3$(0 1 0) and (1 0 0) surfaces, simulated by clusters Mo$_2$O$_{11}$H$_{10}$ and Mo$_3$O$_{12}$H$_6$, respec-

### Table 1

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<th>Atomic charges and bond orders for Mo$<em>7$O$</em>{30}$H$_{18}$, Mo$<em>5$O$</em>{19}$H$_8$, and Mo$<em>3$O$</em>{12}$H$_6$ clusters representing the MoO$_3$(0 1 0) surface</th>
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The charges $q$ and bond orders $P$ involving molybdenum refer to the two inequivalent centers of the clusters, Mo(1) and Mo(2).

Therefore, the results of the largest cluster should be somewhat more reliable. Altogether, the charge differences between equivalent atoms of the different clusters are always small considering the uncertainties involved in calculating atom charges from populations which demonstrates that the surface cluster approach is reasonable for the present system.
tively, are shown in Fig. 4. The calculations yield $V(z)$ always negative for typical molecular reaction distances above the (0 1 0) surface, see Fig. 4(a). This reflects the negative charge accumulation above the surface oxygens. In addition, there is a broad negative minimum above the terminal oxygen sites O(a) which suggests that electrophilic adparticles, like H$^+$ resulting from surface reactions, will be attracted preferentially at these sites and may form local surface bonds. On the (1 0 0) surface there are also electrostatic potential minima above the terminal O(a) centers whereas $V(z)$ is found always positive above Mo centers at the (1 0 0) surface. This suggests that direct bonding between H$^+$ and Mo ions (resulting eventually in bronze type material) will not occur. Moreover, the presence of bare metal centers at the MoO$_3$(1 0 0) surface leads to electrostatic potentials above O(b$'$) and O(d$_1$) oxygen neighbors which are less negative than those above the bridging oxygens on the (0 1 0) surface, see Fig. 4(b).

### 3.2. Hydrogen adsorption and OH desorption at MoO$_3$(0 1 0)

Bond formation of hydrogen or H$^+$ approaching MoO$_3$(0 1 0) and (1 0 0) at different surface oxygen sites is studied in total energy calculations on respective (Mo$_n$O$_m$H$_k$) clusters [6]. The potential energy curves $E_{tot}(z)$ for H$^+$ adsorption and subsequent OH desorption at the MoO$_3$(0 1 0) and (1 0 0) surfaces are shown in Figs. 5 and 6, respectively. Here $z(O-H)$ denotes the perpendicular distance between the adsorbing H$^+$ and the respective surface oxygen center while $z$(Sub-OH) represents the shift of the abstracted oxygen center of the desorbing OH from its initial surface position. The results of Fig. 5 are obtained with the electron exchange and correlation interaction accounted for by both the local spin density (LSDA) and the generalized gradient (GGA) approximation. Comparison shows a very close similarity between the
two approaches. Thus, the curves of Fig. 6 are based on the LSDA scheme only.

Hydrogen is found to stabilize at all sites on both surfaces at similar distances $d(O-H)$ close to the distance of free hydroxyl. Its rather strong binding, resulting in a local OH group, is determined by covalent as well as ionic contributions. The participation of ionic contributions to O–H binding is confirmed by the electrostatic surface potential. Further, a comparison of Fig. 4(a) and (b) with Figs. 5 and 6 shows that terminal oxygen sites O(a) which lead to strongest binding are also sites where the electrostatic potential of the substrate cluster assumes its (negative) minimum. If, after stabilization of hydrogen at the surface, the local OH group is removed at different oxygen sites (simulated by total energy calculations on respective (Mo$_n$O$_{m-1}$H$_k$+OH) clusters), the calculations reveal a qualitative difference between the different sites. The cluster total energy for the removal of OH involving the terminal oxygen O(a) exhibits a sharp minimum near the surface followed by a steep rise. This shows that the respective surface OH group is strongly bound to the underlying Mo metal center. In contrast, the total energy curves for OH removal involving bridging oxygens O(b,b') are rather shallow up to 1.5 Bohr above the surface suggesting that respective surface OH groups can become quite mobile. These hydroxyls are therefore much more available for subsequent desorption and reaction than terminal OH groups. So the present model calculations may suggest that bridging oxygen sites rather than terminal ones participate in the selective oxidation of hydrocarbons at MoO$_3$ surfaces [11]. It should be noted that the present model calculations assume a rigid substrate for OH desorption. Thus, possible surface relaxation and reconstruction (requiring an evaluation of rather complex multi-dimensional interaction potentials) are neglected and the present models can account only for initial desorption steps. However, these initial steps are believed to determine possible reaction channels which will also occur in a full treatment of the surface systems.

3.3. Allyl oxidation at MoO$_3$(0 1 0) and (1 0 0)

The oxidation of allyl (CH$_2$–CH–CH$_2$) to acrolein (CH$_2$=CH–CHO) at MoO$_3$ substrates has been found experimentally to occur for (0 1 0) but not for (1 0 0) oriented surfaces [1]. It can be described by a concerted reaction where a hydrogen is removed from the allyl while a substrate oxygen is incorporated to yield acrolein. In a first attempt to study this reaction on a microscopic basis, and in particular, to understand its surface specificity, cluster calculations have been performed for the adsorption and surface binding of allyl at the MoO$_3$(0 1 0) and (1 0 0) surfaces [7]. Here the local substrate environment near the surface is simulated by saturated Mo$_2$O$_{11}$H$_{10}$ (for MoO$_3$(0 1 0), see Fig. 2(d)) and Mo$_3$O$_{12}$H$_6$ clusters (for MoO$_3$(1 0 0), see Fig. 3(d)) to which the C$_3$H$_5$ adsorbate species is added. While the adsorbate geometry is taken from that of the free radical, the adsorbate–substrate distance is optimized in DFT total energy calculations for stabilization above different surface oxygen sites and for different adsorbate orientations, shown in Figs. 7 and 8.

The allyl binding energies, obtained for the above clusters, are plotted as functions of the adsorbate–substrate separation $z$ at the MoO$_3$(0 1 0) and (1 0 0) surfaces in Figs. 9 and 10. Here $z$ refers to a perpendicular distance as indicated by arrows in the insets of the figures showing the different clusters, see also Figs. 7 and 8. For both surface orientations the calculations yield a clearly preferred adsorbate geometry. On the MoO$_3$(0 1 0) surface the allyl binds most strongly above an asymmetric bridging oxygen site, O(b), with the adsorbate plane perpendicular to the
surface and one peripheral hydrogen, H(1), pointing towards the O(b) center, see geometry (1) of Fig. 7. On MoO\textsubscript{3}(1 0 0) the preferred adsorbate geometry is described by (3) of Fig. 8, where the allyl stabilizes above a bare metal center (these centers do not exist on MoO\textsubscript{3}(0 1 0)) with its adsorbate plane parallel to the surface.

The differences \( \Delta P \) in selected intra-adsorbate, intra-substrate and adsorbate–substrate bond orders in the surface clusters resulting from allyl adsorption are listed in Tables 3 and 4. The numerical results (negative \( \Delta P \) values refer to adsorbate induced bond weakening, positive values to strengthening) are given for the energetically preferred adsorbate geometries, (1) for MoO\textsubscript{3}(0 1 0) (Table 3) and (3) for the MoO\textsubscript{3}(1 0 0) (Table 4), at equilibrium distance and also for an adsorbate–substrate distance decreased by 1 Bohr with respect to equilibrium. Here a comparison for different distances can give useful information about adsorbate–substrate bond formation when the adsorbate is approaching the surface. It is obvious from Figs. 9 and 10 that the calculations yield an overall rather weak interaction of the allyl adsorbate with both MoO\textsubscript{3} surfaces together with unusually large equilibrium distances. This is most likely a consequence of the fixed adsorbate and substrate geometry used for the present models. Surface relaxation combined with adsorbate distortion may change the energetics and could even lead to completely new adsorption/reaction channels. Calculations along these lines require much larger clusters compared to those used in the present work and are out of reach.
with present computational schemes. However, the present restricted models reveal already interesting details for the site dependent surface binding in the initial adsorption steps which can suggest possible mechanisms for the allyl to acrolein oxidation process [7].

On the MoO₃(0 1 0) surface the adsorbate becomes most stable above an asymmetric bridging oxygen
site, O(b), where, however, due to the size of the adsorbate the two neighboring terminal oxygen sites, denoted O(a1) and O(a2), are also affected by the surface interaction, see geometry (1) Fig. 7. In this geometry, bond order analyses yield a slight weakening of all intra-allyl bonds accompanied by the formation of new partial C(1)–O(a1), H(1)–O(a1), and H(1)–O(b) bonds. A comparison of the results of the second and third column of Table 3 show that when allyl approaches the MoO3(0 1 0) surface all allyl-substrate bonds are strengthened. Further, the two C–C bonds in the adsorbate become inequivalent with the C(1)–C(2) bond weakening and the C(2)–C(3) bond strengthening with decreasing adsorbate–substrate distance.

Similar conclusions about bond formation at the surface can be drawn from the charge density difference map shown in Fig. 11. Here electron density differences, shown as shaded contour plots, are obtained from a comparison of the adsorbate cluster result with that of a superposition of the isolated adsorbate and substrate cluster. Thus, Fig. 11 quantifies the electron charge rearrangement due to interaction between the allyl and the MoO3 substrate. Fig. 11(a) shows rather clearly a weak bond formation (accumulation of charge, positive charge density differences shown by full contour lines) between the allyl and the MoO3(0 1 0) surface in addition to the asymmetry in the two C–C bonds in the adsorbate. Thus, the present calculations suggest an asymmetric electronic structure of the C3 skeleton of adsorbed allyl which is qualitatively similar to that in acrolein (C(1)–C(2) single bond and C(2)–C(3) double bond) and which can be interpreted as an initial step in the oxidation process.

On the MoO3(1 0 0) surface the allyl binds most strongly above a bare metal center with its adsorbate plane parallel to the surface, see geometry (3) of Fig. 8. As a result, the adsorbate–substrate coupling is dominated by bonds involving adsorbate π and metal orbitals, cp. Table 4 and Fig. 11(b) and (c). Moreover, the calculations for geometry (3) give no indications of O–C surface bond formation and there is no trend towards asymmetric C–C bonds in the adsorbed allyl. Therefore, the results suggest that the allyl oxidation at the MoO3(1 0 0) surface is less likely to happen. Altogether, the present calculations give a rather simple geometric picture of the initial step of allyl to acrolein conversion at MoO3 based on the surface dependent binding character of the adsorbate which can explain the different results between the (0 1 0) and (1 0 0) surfaces found in the experiment.

### 4. Conclusions

The present cluster model calculations [4–7] yield electronic parameters which are well converged with respect to cluster size. Therefore, they can give a clear picture of the electronic structure and binding near the different oxygen sites at the different MoO3 surfaces, (0 1 0) and (1 0 0) oriented. The ionic nature of the MoO3 compound is confirmed by the calculations. Further, populations as well as bond orders show characteristic differences for geometrically inequivalent oxygen surface sites. Terminal molybdenyl oxygens accumulate the smallest negative charge (−0.4) and their interaction with the environment is described by a double bond with the adjacent Mo center. Asym-
metric bridging oxygens are slightly more negative (−0.5) and characterized by a binding scheme similar to that for the molybdenyl oxygens. Symmetric bridging oxygens become most negative (−0.8) and their binding behavior is described by single bonds with the two neighboring Mo centers. The cluster results obtained for geometrically equivalent oxygens at the MoO₃(0 1 0) and (1 0 0) surfaces are found to be very similar and do not seem to depend strongly on the surface geometry. Therefore, the electronic structure at the two ideal surfaces is determined mainly by their detailed atom arrangement and is

Fig. 11. Shaded contour plots of the charge density difference of: (a) Mo₂O₁₁H₁₀-C₃H₅ referring to geometry (1) on MoO₃(0 1 0) for a cut perpendicular to the surface, (b) Mo₂O₁₂H₆-C₃H₅ referring to geometry (3) on MoO₃(1 0 0) for a cut perpendicular to the surface, and (c) Mo₂O₁₂H₆-C₃H₅ as in (b) but for a cut along the allyl plane. Solid (dashed) contours correspond to positive (negative) values, with increments of 0.005 a.u.
not influenced by major charge redistributions due to substrate surface binding.

Binding of hydrogen at the different oxygen sites of MoO$_3$ results always in local surface OH groups and is determined by covalent as well as ionic contributions. Removal of OH from the surface after hydrogen adsorption proceeds differently depending on the surface oxygen site. Here, OH groups involving terminal oxygens O(a) are quite strongly bound to the surface and are thus unlikely to desorb. In contrast, hydroxyls involving bridging oxygens O(b,b') are more loosely coupled to the surface and mobile about the H adsorption minimum. Therefore, they are available for subsequent reactions and may desorb from the surface. So the present model calculations suggest that bridging oxygen sites rather than terminal ones participate in the selective oxidation of hydrocarbons at MoO$_3$ surfaces.

As an example of more complex adsorbate–substrate binding the interaction of allyl with MoO$_3$(0 1 0) and (1 0 0) surfaces is examined [5,7,11]. This interaction is found to be rather weak which may be, to a major extent, due to computational constraints assuming fixed adsorbate and substrate geometry. However, the calculations can give a simple qualitative picture based on the surface dependent binding character which suggests a starting point in the allyl to acrolein oxidation process on the (0 1 0) but not on the (1 0 0) surface. This confirms the experimental findings.

**Acknowledgements**

One of the authors (AM) gratefully acknowledges the Fellowship for Young Scientists from the Foundation for Polish Science. This work was funded in parts by the Deutsche Forschungsgemeinschaft and Fonds der Chemischen Industrie.

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